

## The Crystal Structure of *N*-Methylrheoagenine Iodide

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The structure and stereochemistry of *N*-methylrheoagenine iodide have been determined by three-dimensional Fourier and least-squares methods using visually estimated intensity data. The crystals have orthorhombic symmetry, space group  $C222_1$ , with  $a = 8.54$ ,  $b = 18.52$ ,  $c = 26.19$  Å,  $Z = 8$ . The two 1,3-benzodioxole systems, which are linked by *cis*-fused azacycloheptene and dihydropyran rings, are nearly mutually perpendicular. The hydroxyl group and the hydrogen atoms at the junction of the two heterocyclic rings lie on the same side of the dihydropyran ring. Infinite chains of alternating iodide ions and organic cations are arranged parallel to the [110] direction.

### Introduction

The structures of the alkaloid rheoagenine and the closely related alkaloid rheoadine have been under investigation by chemical methods for several decades. A structure (Fig. 1) was proposed for rheoagenine by Santavy, Kaul, Hruban, Dolejs, Hanus, Blaha & Cross (1965). However, the authors noted that the product of one of the oxidation reactions was still not explained, and since the proposed structure represented a new type of alkaloid skeleton, an X-ray analysis of the methiodide derivative was undertaken. The structure found by X-ray diffraction agrees with the structure proposed on chemical grounds, and also establishes the stereochemistry of the molecule.

### Experimental

The crystals used for the structure determination were kindly provided by Dr Danuta Rozwadowska, then of Carleton University. Crystal data are as follows:

*N*-Methylrheoagenine iodide,  $[C_{21}H_{22}NO_6]^+I^-$   
M.W. 511.3

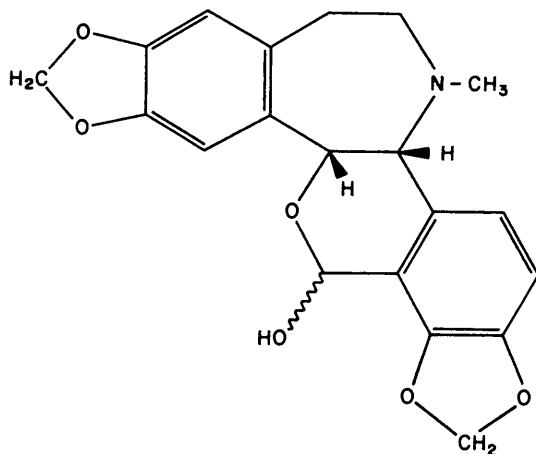


Fig. 1. Structure proposed for rheoagenine.

Orthorhombic, space group  $C222_1$

$a = 8.54 \pm 0.02$ ,  $b = 18.52 \pm 0.04$ ,  $c = 26.19 \pm 0.06$  Å.

$D_m$  (by flotation) =  $1.65$  g.cm $^{-3}$ ,  $Z = 8$ ,  $D_x = 1.64$  g.cm $^{-3}$ .

$\mu$  (Cu  $K\alpha$ ) =  $128.1$  cm $^{-1}$ .

Unit-cell dimensions were determined from precession photographs obtained with Cu  $K\alpha$  radiation ( $\lambda = 1.5418$  Å). Systematic absences  $hkl$  when  $h+k$  is odd and  $00l$  when  $l$  is odd were deduced from Weissenberg and precession photographs, and led uniquely to space group  $C222_1$ .

Intensity data were collected from equi-inclination multiple-film Weissenberg photographs obtained with Cu  $K\alpha$  radiation for the  $0kl \cdot \cdot \cdot 7kl$ ,  $h0l \cdot \cdot \cdot h4l$ , and  $hk0$  layers. The crystals deteriorated noticeably with exposure to X-rays, and consequently fresh crystals were mounted at intervals during the data collection. Altogether, seven different crystals were used, each of which was ground to an approximately spherical shape in order to simplify the application of absorption corrections.

Intensities were estimated by visual comparison with a calibrated intensity scale, and corrected for Lorentz and polarization factors. Absorption corrections appropriate for spherical crystals (*International Tables for X-ray Crystallography*, 1959) were applied to the various sets of data collected with different crystals. The  $F^2$  values from various layers were placed on a common relative scale using the Hamilton, Rollett & Sparks (1965) method. Of the 2643 independent reflexions accessible with Cu  $K\alpha$  radiation, 2137 were observed, and another 316 reflexions were recorded as having intensities too weak to be measured. The latter were excluded from the structure determination and refinement, but were included in the final structure factor calculation; the value given in place of  $|F_o|$  was two thirds of the minimum observable amplitude.

### Structure determination and refinement

The coordinates of the iodide ion were deduced from two-dimensional Patterson syntheses, and the first

three-dimensional heavy-atom-phased Fourier synthesis revealed all twenty-eight atoms of the organic cation. A subsequent structure-factor calculation in which all non-hydrogen atoms were included gave an *R* value ( $= \sum |K|F_o - |F_c| / \sum |K|F_o$  for all observed reflexions) of 0.24.

The structure was refined by block-diagonal least-squares calculations. The quantity minimized was  $\sum w(K|F_o - |F_c|)^2$ , and the form of the weighting scheme used throughout the refinement was

$$w = 1/[1 + \{(K|F_o - g)/f\}^2].$$

Table 1. Fractional coordinates, vibration tensor components ( $\text{\AA}^2$ ) for the expression  $T = \exp[-2\pi^2(U_{11}a^*h^2 + \dots + 2U_{23}b^*c^*kl + \dots)]$ , and their e.s.d.'s (all quantities  $\times 10^4$ )

	<i>x</i>	<i>y</i>	<i>z</i>	<i>U</i> <sub>11</sub>	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	2 <i>U</i> <sub>23</sub>	2 <i>U</i> <sub>13</sub>	2 <i>U</i> <sub>12</sub>
I	1808 (1)	2254 (1)	1211 (1)	725 (6)	528 (4)	907 (7)	-337 (10)	-303 (11)	196 (10)
N	6966 (15)	3340 (7)	1193 (5)	541 (65)	643 (61)	719 (63)	-5 (111)	86 (129)	-5 (106)
O(1)	8101 (13)	4891 (5)	1050 (3)	691 (55)	544 (43)	483 (38)	39 (66)	150 (86)	138 (95)
O(2)	8967 (17)	5758 (7)	1605 (4)	861 (75)	735 (64)	671 (56)	312 (98)	51 (114)	-93 (123)
O(3)	13932 (15)	4198 (7)	48 (4)	674 (57)	1166 (87)	417 (42)	25 (103)	139 (92)	-146 (135)
O(4)	12647 (14)	3804 (8)	-656 (4)	678 (64)	1079 (87)	542 (49)	-460 (108)	116 (96)	210 (134)
O(5)	6956 (16)	5746 (5)	2551 (4)	857 (77)	679 (52)	542 (45)	-74 (84)	263 (110)	66 (119)
O(6)	6715 (17)	5037 (7)	3268 (4)	995 (81)	834 (63)	521 (45)	-66 (91)	405 (119)	-296 (151)
C(1)	7213 (21)	3116 (11)	620 (8)	588 (89)	879 (111)	893 (104)	-499 (182)	71 (162)	-231 (169)
C(2)	7571 (20)	3680 (13)	250 (6)	480 (77)	1265 (148)	632 (76)	-632 (184)	-210 (137)	309 (197)
C(3)	9210 (17)	3891 (8)	197 (5)	522 (72)	627 (71)	497 (59)	22 (106)	7 (109)	127 (124)
C(4)	9992 (18)	4154 (9)	627 (5)	520 (71)	759 (82)	438 (56)	-69 (116)	-122 (108)	30 (137)
C(5)	9203 (15)	4307 (7)	1139 (5)	446 (64)	530 (59)	549 (66)	32 (101)	-334 (108)	113 (102)
C(6)	8436 (17)	3703 (7)	1424 (5)	519 (72)	533 (61)	546 (59)	76 (98)	-80 (112)	-68 (119)
C(7)	7973 (19)	4029 (7)	1943 (5)	668 (89)	473 (57)	539 (59)	179 (96)	-218 (127)	1 (123)
C(8)	7674 (18)	4777 (8)	1967 (5)	565 (74)	619 (75)	487 (59)	-17 (108)	2 (112)	113 (124)
C(9)	7798 (19)	5274 (7)	1509 (4)	708 (89)	533 (63)	410 (52)	22 (93)	71 (112)	164 (123)
C(10)	11602 (18)	4273 (7)	6142 (5)	561 (74)	565 (64)	483 (56)	-192 (96)	187 (116)	67 (123)
C(11)	12412 (17)	4125 (8)	159 (5)	473 (66)	648 (73)	542 (63)	160 (113)	13 (110)	280 (127)
C(12)	11570 (25)	3897 (9)	-264 (5)	1014 (130)	706 (83)	427 (59)	98 (111)	-90 (151)	416 (188)
C(13)	9995 (23)	3759 (10)	-273 (6)	735 (99)	805 (92)	553 (73)	-147 (140)	126 (142)	225 (177)
C(14)	14015 (29)	4074 (13)	-481 (6)	1002 (133)	1100 (136)	525 (73)	-312 (170)	105 (179)	47 (248)
C(15)	7293 (18)	5036 (8)	2422 (5)	604 (80)	582 (63)	563 (70)	143 (116)	96 (121)	45 (131)
C(16)	7257 (19)	4626 (9)	2871 (4)	676 (86)	818 (87)	327 (45)	125 (108)	32 (107)	-80 (150)
C(17)	7568 (21)	3912 (9)	2865 (5)	781 (96)	834 (96)	330 (49)	238 (113)	78 (118)	99 (166)
C(18)	7968 (18)	3610 (8)	2382 (5)	519 (84)	737 (78)	553 (70)	253 (116)	-19 (120)	161 (137)
C(19)	6909 (23)	5783 (8)	3090 (6)	706 (95)	659 (78)	702 (76)	-7 (128)	203 (167)	-230 (169)
C(20)	6628 (30)	2660 (11)	1462 (11)	906 (134)	688 (106)	1442 (177)	-140 (209)	-62 (256)	-672 (216)
C(21)	5500 (18)	3784 (11)	1196 (7)	460 (74)	935 (104)	796 (90)	34 (179)	68 (151)	58 (156)

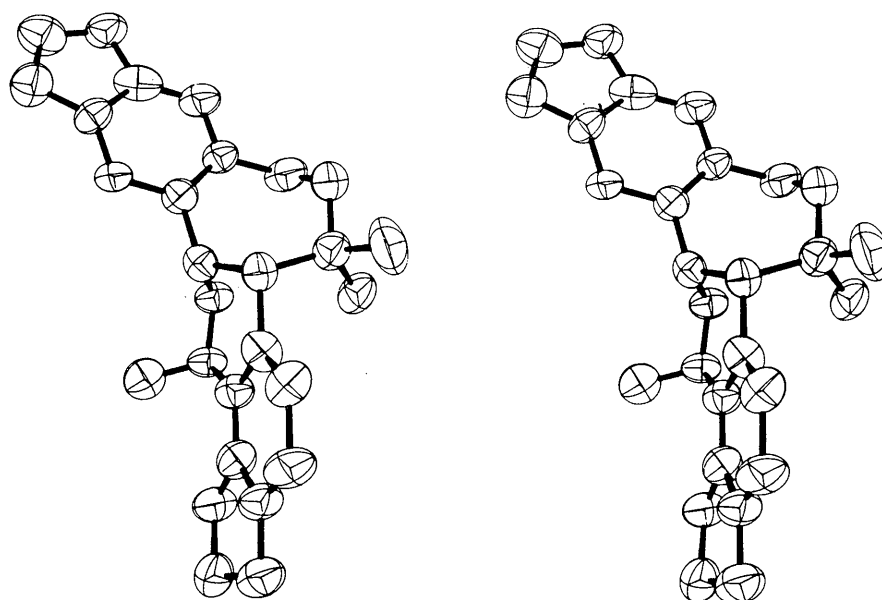


Fig. 2. Stereoscopic view of the organic cation in *N*-methylrhoeagenine iodide. The ellipsoids are scaled to include 50% probability.

Table 2. Observed and calculated structure amplitudes (x 10)

\* Indicates an unobserved reflexion, with 2/3 F\_min given in place of F\_0.

Table with multiple columns for h, k, l indices and corresponding observed and calculated structure amplitudes. The table is organized into several vertical sections, each with its own set of indices and amplitude values. Asterisks indicate unobserved reflexions.

Table 2 (cont.)

L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC	L	FO	FC															
16 252 282			17 188 170			18 243 197			19 156 467			20 246 236			21 195 190			22 249 238			23 195 190			24 611 114			25 165 156			26 165 156			27 182 185			28 185 185			29 137 125			30 137 125		
16 252 282			17 188 170			18 243 197			19 156 467			20 246 236			21 195 190			22 249 238			23 195 190			24 611 114			25 165 156			26 165 156			27 182 185			28 185 185			29 137 125			30 137 125		
16 252 282			17 188 170			18 243 197			19 156 467			20 246 236			21 195 190			22 249 238			23 195 190			24 611 114			25 165 156			26 165 156			27 182 185			28 185 185			29 137 125			30 137 125		

Table 3. Magnitudes  $B$  and direction cosines relative to  $a$ ,  $b$ , and  $c$ , of the principal axes of the thermal ellipsoids

	$B_1$	$B_2$	$B_3$	$D_{1a}$	$D_{1b}$	$D_{1c}$	$D_{2a}$	$D_{2b}$	$D_{2c}$	$D_{3a}$	$D_{3b}$	$D_{3c}$
I	5.05	8.40	3.61	-0.867	-0.012	-0.498	-0.465	-0.342	0.816	-0.180	0.940	0.291
N	4.20	5.76	5.08	0.975	0.016	-0.223	0.224	-0.034	0.974	0.008	-0.999	-0.037
O(1)	5.85	4.09	3.62	0.896	0.341	0.285	0.312	-0.939	0.141	-0.316	0.037	0.948
O(2)	6.67	6.99	4.23	-0.642	-0.523	-0.561	-0.752	0.572	0.327	0.151	0.632	-0.760
O(3)	5.38	9.29	3.14	-0.956	-0.137	-0.258	-0.143	0.990	0.003	-0.255	-0.040	0.966
O(4)	5.52	9.29	3.35	-0.925	0.026	-0.379	-0.159	0.933	-0.323	-0.345	0.359	0.867
O(5)	7.15	5.44	3.82	0.940	0.082	-0.331	-0.001	-0.970	0.243	-0.341	0.229	0.912
O(6)	9.03	6.02	3.52	0.848	-0.438	0.299	-0.391	-0.897	-0.205	-0.358	-0.057	0.932
C(1)	4.03	9.12	5.28	0.857	0.459	0.235	0.186	-0.700	0.690	0.481	-0.547	-0.685
C(2)	4.11	11.30	3.37	0.431	-0.426	-0.795	0.189	0.904	-0.382	0.882	-0.015	0.471
C(3)	3.92	5.20	3.88	0.314	-0.073	-0.947	0.424	0.903	0.071	0.849	-0.423	0.315
C(4)	4.33	6.04	3.19	-0.885	0.135	0.446	0.090	0.989	-0.121	0.457	0.067	0.887
C(5)	4.27	5.33	2.44	-0.105	-0.959	-0.265	-0.606	-0.150	0.781	0.789	-0.242	0.565
C(6)	3.96	4.81	3.85	-0.009	-0.765	0.644	-0.505	0.560	0.657	0.863	0.317	0.393
C(7)	5.83	4.33	3.11	0.828	-0.178	-0.531	-0.500	-0.664	-0.556	0.254	-0.726	0.639
C(8)	4.18	5.17	3.83	-0.843	0.523	0.126	0.531	0.846	-0.040	-0.086	0.101	0.991
C(9)	5.88	3.96	3.20	0.925	0.364	0.111	0.363	-0.931	0.029	-0.114	-0.013	0.993
C(10)	4.71	5.07	2.93	-0.735	-0.677	-0.033	0.505	-0.576	0.643	-0.455	0.455	0.765
C(11)	4.14	5.93	3.06	0.366	0.175	-0.914	0.429	0.840	0.333	0.826	-0.514	0.232
C(12)	8.83	4.91	3.20	0.894	0.447	-0.026	0.417	-0.853	-0.312	0.162	-0.268	0.950
C(13)	5.59	7.02	3.91	-0.727	0.458	-0.511	0.568	0.820	-0.074	-0.386	0.344	0.856
C(14)	7.96	9.00	3.78	-0.992	0.043	-0.122	0.072	0.968	-0.239	-0.108	0.246	0.963
C(15)	4.55	5.35	3.91	-0.803	0.565	0.191	0.564	0.575	0.593	0.223	0.582	-0.782
C(16)	5.27	6.60	2.51	-0.971	-0.223	-0.086	-0.232	0.966	0.112	-0.058	-0.129	0.990
C(17)	5.95	7.03	2.38	-0.884	0.465	0.049	0.464	0.859	0.215	-0.058	-0.213	0.975
C(18)	4.27	6.45	3.56	-0.744	-0.119	0.658	0.225	0.882	0.414	0.629	-0.456	0.629
C(19)	5.37	6.70	4.25	0.142	-0.636	-0.758	0.729	-0.451	0.515	0.669	0.626	-0.400
C(20)	9.08	11.44	3.46	-0.811	0.583	0.053	-0.001	-0.091	0.996	0.585	0.808	0.074
C(21)	3.59	7.42	6.29	0.994	-0.056	-0.096	0.069	0.988	0.135	-0.088	0.141	-0.986

The parameters  $f$  and  $g$  were re-evaluated several times so as to make  $\langle w(K|F_o| - |F_c|)^2 \rangle$  essentially independent of  $|F_o|$ . Atomic scattering factor values were those of Hanson, Herman, Lea & Skillman (1964) for the carbon, nitrogen and oxygen atoms and of Cromer & Waber (1965) for the iodide ion. The real part of the anomalous dispersion correction given by Cromer (1965) was applied to the  $I^-$  scattering curve.

Three cycles of least-squares calculations using isotropic temperature factors for all atoms, followed by three cycles in which the iodide atom was refined anisotropically and the other atoms isotropically, resulted in an  $R$  index of 0.12. After several further cycles in which all atoms were refined anisotropically, it became clear that agreement between observed and calculated structure factors was generally worse for reflexions with  $\sin^2\theta \leq 0.200$  than for the rest of the data. It seems probable that inaccuracy in the values of  $\mu R$  assumed

for the absorption corrections, and, especially, deviation of the crystals from strictly spherical shape, may be the main factors responsible for the low-angle discrepancies. (Crystal deterioration can be expected to affect all the data adversely.) Because of the apparent unreliability of the low-angle data, reflexions with  $\sin^2\theta \leq 0.200$  were excluded from further refinement. After a few more least-squares cycles, refinement was terminated; the average parameter shifts were then less than one tenth of their respective estimated standard deviations, and the largest shift was  $0.8\sigma$ . The final  $R$  index for the 1944 reflexions used in the last cycles of refinement was 0.092, while for all 2137 observed reflexions it was 0.113.

The estimated standard deviations of the atomic parameters were calculated from the diagonal elements of the inverse to the normal equations matrix. The e.s.d.'s of the bond lengths and angles which were derived from these have been increased by 50% to allow for omission of unobserved reflexions and interactions between atoms.

A difference map computed with the final structure factors showed a number of peaks which could be interpreted as hydrogen atom sites, but only nine of these, less than half the total number of hydrogen atoms, were significantly higher than the background level. Consequently no attempt was made to pursue the refinement including hydrogen atoms.

## Results and discussion

The molecular structure and configuration may be seen in Fig. 2, which is a stereoscopic view of the organic cation, prepared with Johnson's (1965) ORTEP program.

Final coordinates and anisotropic thermal parameters with corresponding e.s.d.'s are given in Table 1. Observed and calculated structure amplitudes based on these parameters are listed in Table 2. Magnitudes and direction cosines of the principal axes of the individual ellipsoids are given in Table 3.

An analysis of the thermal parameters by Cruickshank's (1956) method, assuming the 28 non-hydrogen atoms of the organic cation to act as a rigid body, indicated that the main rotational vibration was about an axis roughly parallel to the longest dimension of the cation, and that the translational vibration of the ion was nearly isotropic. However, the discrepancies between the observed  $U_{ij}$ 's and the  $U_{ij}$ 's calculated from the  $T_{ij}$  and  $\omega_{ij}$  tensors were much larger than the e.s.d.'s of the observed  $U_{ij}$  values. Since the rigid-body approximation does not appear to be appropriate in this structure, no corrections to the bond lengths were calculated from the  $\omega_{ij}$  tensor values.

Corrections were, however, made to the C(9)–O(2), N–C(20) and N–C(21) bond lengths, using the formulae for 'riding' motion given by Busing & Levy (1964). The corrections for these three bonds amounted to 0.019, 0.033 and 0.010 Å respectively.

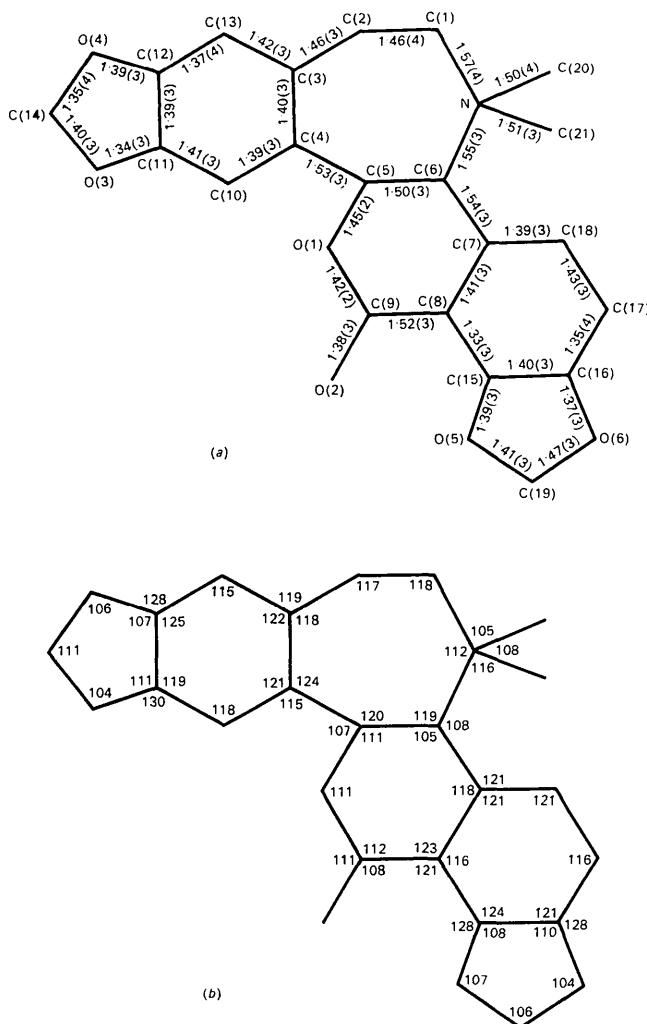


Fig. 3. Bond lengths (Å) and their e.s.d.'s ( $\times 10^2$ ), and bond angles (°). E.s.d.'s for the angles range from 1.5 to 2.5°.

**Bond lengths and angles**

The numbering scheme and bond lengths with their e.s.d.'s are shown in Fig. 3(a), and bond angles are given in Fig. 3(b). The two angles not shown in the Figure for simplicity are C(1)–N–C(21), 105°; C(6)–N–C(20), 110°. The values given in Fig. 3(a) for the C(9)–O(2), N–C(20) and N–C(21) bonds are the corrected bond lengths.

The four N<sup>+</sup>–C *sp*<sup>3</sup> bonds have an average length of 1.53 Å, and although the two bonds in the seven-membered ring appear longer than the other two, the difference is not significant. This average may be compared with the value 1.52 ± 0.03 Å quoted by Hamilton, Hamor, Robertson & Sim (1962) as characteristic for the N<sup>+</sup>–C *sp*<sup>3</sup> bond length in a number of alkaloids.

As would be expected, the average of the seven C *sp*<sup>3</sup>–O bonds, 1.41 Å, is longer than the average C *sp*<sup>2</sup>–O bond length, 1.37 Å, and each is in satisfactory agreement with respective standard values 1.426 ± 0.005 and 1.37 ± 0.02 Å (Sutton, 1965) or 1.425 and 1.381 Å (Hall & Ahmed, 1968). The two C *sp*<sup>3</sup>–C *sp*<sup>3</sup> bonds appear shorter than the usual value of 1.537 Å (Sutton, 1965), but their deviations are not significant. The ave-

rage C *sp*<sup>2</sup>–C *sp*<sup>3</sup> bond length, 1.51 Å, is also in agreement with the characteristic value 1.510 ± 0.005 Å (Sutton, 1965).

For the twelve aromatic C *sp*<sup>2</sup>–C *sp*<sup>2</sup> bonds, the average length is 1.39 Å. The deviation of one bond, C(8)–C(15), from the average is in the range of possible significance (1% < *P* < 5%) using the *t* test, but since there is no obvious reason for its shortness, this may indicate that the e.s.d.'s are slightly underestimated.

The coordination of the positively charged nitrogen atom is approximately tetrahedral, but the angle C(6)–N–C(21) is significantly enlarged from 109.5° because of repulsion between C(21) and O(1). Even with this angular distortion, C(21) is only 3.04 Å from O(1), whereas the normal CH<sub>3</sub>...O van der Waals separation is 3.4 Å (Pauling, 1960).

The angles in the azacycloheptene ring are all larger than 109.5°; those at C(3) and C(4) are close to 120° as would be expected because of the fused aromatic ring. The average of the other five angles in the ring, 117°, is similar to the values found for cycloheptane ring angles in various compounds; for example: isoclovene hydrochloride (Clunie & Robertson, 1961), 116.5°; bromo-

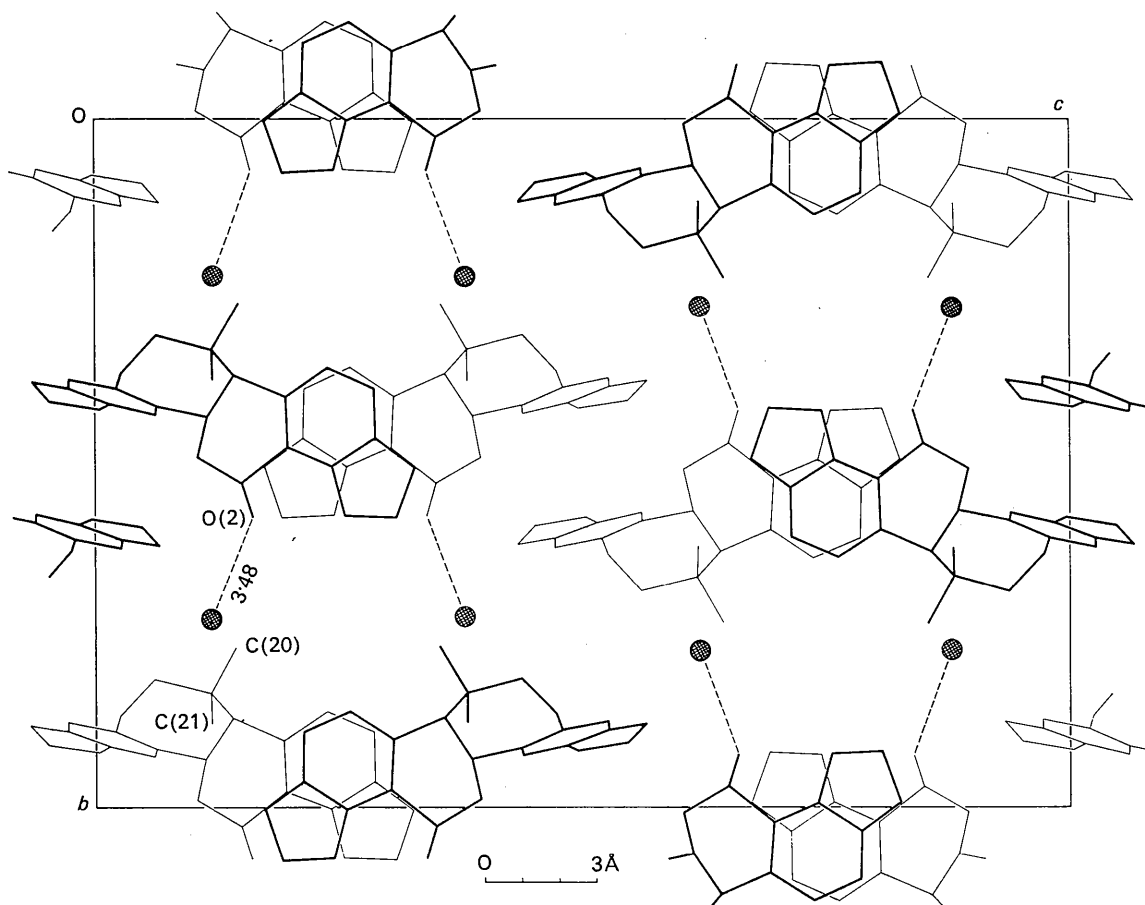


Fig. 4. Projection of the structure on the (100) plane. The hatched circles represent iodide ions, and the proposed hydrogen bonding between the iodide ion and hydroxyl oxygen atom is indicated by dashed lines.

geigerin acetate (Hamilton, McPhail & Sim, 1962),  $116^\circ$ ; and 2-bromodihydroisophoto- $\alpha$ -santonin lactone acetate (Asher & Sim, 1965),  $115^\circ$ .

Angles in the benzodioxole groups are similar, within the limits of accuracy of this analysis, to those found by Hall & Ahmed (1968) in the corresponding parts of protopine.

#### Planarity and conformation of the rings

Both aromatic rings are planar within experimental error, since  $\chi^2 = 4.9$  and  $3.7$  for the planes including C(3) and C(7) respectively, and for  $\nu = 3$ ,  $P > 5\%$ . Equations of these two planes are respectively

$$-0.1615X + 0.9432Y - 0.2904Z - 5.3912 = 0$$

$$\text{and } -0.9670X - 0.1962Y - 0.1632Z + 8.8923 = 0$$

where the coefficients of  $X$ ,  $Y$ , and  $Z$  are direction cosines with respect to the cell axes  $a$ ,  $b$  and  $c$ . The dioxole ring formed by O(3), O(4), C(11), C(12) and C(14) is slightly envelope shaped, with O(3), O(4), C(11) and C(12) coplanar within  $0.012 \text{ \AA}$  and C(14) deviating from this plane by  $0.16 \text{ \AA}$ . The other dioxole ring is slightly more puckered, with O(5), O(6), C(15) and C(16) deviating from planarity by amounts up to  $0.044 \text{ \AA}$ , and C(19) lying  $0.30 \text{ \AA}$  from this approximate plane. The two 1,3-benzodioxole systems are nearly mutually perpendicular; the dihedral angle between the two aromatic rings is  $88.9^\circ$ .

The dihydropyran ring is 'sofa'-shaped (Philbin & Wheeler, 1958). As required by the adjacent aromatic ring C(6), C(7), C(8) and C(9) are coplanar, and O(1) is nearly in this plane although its deviation of  $0.12 \text{ \AA}$  is highly significant ( $P < 0.1\%$ ). C(5), however, deviates by  $0.66 \text{ \AA}$  from the plane. The deviation of the hydroxyl oxygen atom, O(2) is  $1.16 \text{ \AA}$ , on the same side of the plane as C(5). This orientation of the hydroxyl group agrees with the recent stereochemical assignment, from nuclear magnetic resonance evidence, by Shamma, Weiss, Pfeifer & Döhnert (1968) for this group in rhoegenine. The dihydropyran and azacycloheptene rings are *cis*-fused, and the ring junction is characterized by dihedral angles of  $63.0^\circ$  for O(1)-C(5)-C(6)-C(7) and  $66.4^\circ$  for C(4)-C(5)-C(6)-N.

The azacycloheptene ring is in a slightly distorted chair form. Despite the fused aromatic ring, C(2), C(3), C(4) and C(5) deviate significantly from planarity (by amounts up to  $0.027 \text{ \AA}$ ). The dihedral angle between the mean plane through these four atoms and the plane defined by C(1), N and C(6) is  $5.0^\circ$ . C(1), C(2), C(5) and C(6) are very significantly non-planar, with deviations up to  $0.077 \text{ \AA}$  from the mean plane through these four atoms. These distortions from normal chair conformation help to increase slightly the distance between C(21) and O(1) and thereby reduce their mutual repulsion.

#### Packing arrangement

Each iodide ion is equidistant,  $4.24 \text{ \AA}$ , from the two methyl groups of one adjacent cation, and is  $3.48 \text{ \AA}$

from the hydroxyl oxygen atom of another cation. The  $I \cdots C(20)$  and  $I \cdots C(21)$  distances are slightly larger than the normal  $I \cdots CH_3$  van der Waals separation,  $4.15 \text{ \AA}$  (Pauling, 1960). However, the equality of the distances suggests that the packing arrangement may be partially governed by Coulomb attraction between the iodide ion and the positively charged but inaccessible nitrogen atom.

The  $I \cdots O$  contact mentioned above is slightly shorter than the usual separation,  $3.55 \text{ \AA}$  (Pauling, 1960) for this type of van der Waals interaction. Since the standard deviation of the  $I \cdots O$  distance is only  $0.02 \text{ \AA}$ , the shortening is considered significant, and this suggests that a weak hydrogen bond may exist between O(2) and  $I^-$ . The value of the  $C(9) - O(2) \cdots I^-$  angle,  $94.7^\circ$ , is also compatible with the existence of such a hydrogen bond.

Contact distances less than  $3.5 \text{ \AA}$  are listed in Table 4. With the exception of the  $I^- \cdots O(2)$  and  $C(11) \cdots C(11)$  distances, none is shorter than the corresponding van der Waals contact distance; the  $C(11) \cdots C(11)$  distance with an e.s.d. of  $0.03 \text{ \AA}$ , is not significantly shorter than the usual van der Waals distance.

Table 4. Intermolecular distances less than  $3.5 \text{ \AA}$

O(2) <sup>i</sup>	I <sup>ii</sup>	3.48 Å	
C(2) <sup>i</sup>	O(3) <sup>iii</sup>	3.30	
C(21) <sup>i</sup>	O(3) <sup>iii</sup>	3.38	
C(14) <sup>i</sup>	O(3) <sup>iv</sup>	3.40	
O(3) <sup>i</sup>	O(3) <sup>iv</sup>	2.98	
C(11) <sup>i</sup>	O(3) <sup>iv</sup>	3.41	
C(19) <sup>i</sup>	O(4) <sup>v</sup>	3.40	
O(5) <sup>i</sup>	O(5) <sup>vi</sup>	3.35	
C(21) <sup>i</sup>	O(6) <sup>vi</sup>	3.31	
C(11) <sup>i</sup>	C(11) <sup>iv</sup>	3.35	
i	x	y	z
ij	$\frac{1}{2} + x$	$\frac{1}{2} + y$	z
iii	x - 1	y	z
iv	x	1 - y	-z
v	2 - x	1 - y	$\frac{1}{2} + z$
vi	1 - x	y	$\frac{1}{2} - z$

The packing arrangement is illustrated in Fig. 4. The structure can be thought of as consisting of infinite chains, parallel to  $[110]$ , of alternating organic cations and iodide ions. The protruding parts of adjacent chains along the  $c$  axis are closely interleaved, as is clearly shown in the Figure.

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## Structure Cristalline et Moléculaire du (Dinitro-2',4'-phényl)-1-chloro-4-pyrazole

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The crystal structure of 1-(2',4'-dinitrophenyl)-4-chloropyrazole has been determined from three-dimensional data collected on a single-crystal diffractometer with Cu  $K\alpha$  radiation. The crystals are orthorhombic, space group  $Pna2_1$  with  $a = 15.21 \pm 0.02$ ;  $b = 13.59 \pm 0.02$ ;  $c = 5.20 \pm 0.01$  Å;  $Z = 4$ . The structure was refined by least-squares methods and the final  $R$  is 0.051 for 822 reflexions. The only significant difference from the structure of 1-(2',4'-dinitrophenyl)-4-bromopyrazole is in the three bonds lengths C(3)-N(4), C(7)-N(5) and C(10)-N(14). The small displacements from the plane of the phenyl ring are similar in the two compounds.

Dans une précédente communication nous avons donné la structure cristalline du (dinitro-2',4'-phényl)-1-bromo-4-pyrazole (Galigné & Falgueirettes, 1969). La détermination de la structure du (dinitro-2',4'-phényl)-1-chloro-4-pyrazole a été entreprise en vue de préciser certains détails concernant la déformation du cycle benzénique et la géométrie du noyau pyrazole.

La synthèse et l'étude chimique de ces produits ont été réalisées dans le laboratoire de M. le Professeur Jacquier (Elguero, Jacquier & Tien Duc, 1966).

### Partie expérimentale

La préparation des cristaux, les mesures des intensités des réflexions à l'aide d'un diffractomètre à compteur ont été effectuées dans les mêmes conditions que pour le (dinitro-2',4'-phényl)-1-bromo-4-pyrazole. Nous avons utilisé les intensités des réflexions pour lesquelles  $\sigma/N$  était inférieur à 3% [ $\sigma = \text{écart type}$ ,  $N = N(\text{pic}) - N(\text{fond continu})$ ]. Sur un total de 1081 réflexions possib-

les, 822 avaient une intensité répondant à cette condition.

### Données cristallographiques

$a = 15,21 \pm 0,02$  Å

$b = 13,59 \pm 0,02$

$c = 5,20 \pm 0,01$

Groupe spatial  $Pna2_1$

$Z = 4$

$V = 1082$  Å<sup>3</sup>

Densité mesurée  $D_m = 1,63 \pm 0,04$  g.cm<sup>-3</sup>

Densité calculée  $D_c = 1,65$  g.cm<sup>-3</sup>

Coefficient linéaire d'absorption  $\mu_l = 24,68$  cm<sup>-1</sup>

### Détermination de la structure

Les deux structures étant isotopes, les coordonnées et paramètres d'agitation thermique obtenus en fin d'affinement avec agitation thermique isotrope du (dinitro-2',4'-phényl)-1-bromo-4-pyrazole ont été utilisés com-